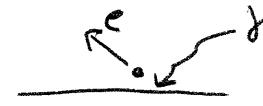


Planck had announced result on Oct 19, 1900, and found  $h = 6.55 \times 10^{-27}$  erg-sec  
 (radiation law since  
 has been beautifully  
 confirmed by better data.) Cf  $h = 6.6262 \times 10^{-27}$  erg-sec  
 - modern value)

But Planck did much more than fit the curve. He announced a derivation of the radiation law on Dec 14, 1900. In what he described as "an act of desperation" he hypothesized that radiation could be emitted and absorbed by the cavity only in quantized form  $\Delta E = n \hbar \omega$

that a new era in physics had begun was not yet appreciated.

### photoelectric effect ("second coming of h")

Lenard 1902 finds electron energy  independent of light intensity

Einstein 1905: an explanation

He radically extends Planck's idea, by proposing that quantization applies not just to interaction of radiation with the cavity, but to radiation itself. Light consists of quanta with

$$E_{\text{quantum}}(\omega) = \hbar \omega$$

("light-quantum hypothesis")

Thus electron energy will be

$$E_{max} = h\nu - \Phi \text{ in "workfunction" of metal}$$

Brilliantly confirmed by Millikan (then at U.Chicago)  
In 1915, he finds  $h = 6.57 \times 10^{-21}$  erg-sec

Why was this radical? Because Einstein was tampering with the greatest success of 19<sup>th</sup> century physics — Maxwell's em theory, which established that light is an electro-magnetic wave. It was one thing to suggest new phenomena in the interaction of radiation with matter (Planck); quite another to suggest that the (Maxwell) theory of free em field is wrong.

Noone accepted this hypothesis

Planck (1913), in letter nominating Einstein for Prussian Academy of Sciences:

"That he may sometimes have missed the target with his speculations, as, for example, in his hypothesis of light-quanta, cannot really be held against him, for it is not possible to introduce really new ideas even in the most exact sciences without sometimes taking a risk."

Millikan (1915), in his paper on the experiments

= Einstein's photoelectric equation ---- appears  
in every case to predict exactly the  
observed results --- yet the  
semi-corpuscular theory by which Einstein  
arrived at his equation seems at  
present to be wholly untenable."

The tide was eventually turned  
by the Compton effect - electron "photon"  
scattering explained by  
relativistic kinematics  
 $E = h\nu$        $\vec{p} = \frac{h}{c}\vec{k}$        $\Rightarrow$   
(idea of photon momentum  
introduced by Einstein in 1916.)       $d\lambda = \frac{h}{mc} (1 - \cos\theta)$

This was 1923. (The acceptance of  
wave-particle duality for photon enabled  
deBroglie to propose same for electron the  
next year.)

The modern derivation of Planck radiation  
law (counting photons) due to Bose  
and Einstein (1924).

It was Einstein, in 1905 paper  
on photoelectric effect, who first  
clearly stated that Planck's radiation  
law created a crisis for  
classical physics. Applying concepts  
of classical statistical mechanics  
(equipartition theorem), he derived

$$u_w = \frac{\omega^2}{\pi^2 c^3} T$$

(the  $T \rightarrow 0$  limit of Planck's formula)  
 Also derived by Rayleigh and Jeans  
 in 1905, and known as the "Rayleigh-Jeans  
 Law". Einstein recognized that  
 this law must fail at high frequency  
 (or  $V/V = \int du_w = \infty$ ),  
 as indeed it does according to Planck's  
 law. Something radically new had  
 entered physics, something boy and  
 classical concepts (Jeans thought to arise because high-  
 frequency modes are not in equilibrium).

### Debye Theory of Phonons

— (sound quanta)

uses similar mathematics to derivation of  
 Planck's law, but applies it to a rather  
 different physical context.

Consider vibrations of a crystal about  
 equilibrium position. For small  
 vibrations, restoring forces  
 are harmonic. Like a system  
 of masses and springs (coupled oscillators)  
 can be expanded in terms of normal modes,  
 which behave like independent uncoupled  
 oscillators — each with characteristic  $\omega$

For a large crystal, we may think of  
 the modes as traveling (sound)

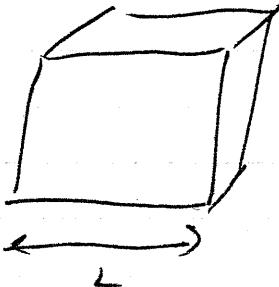
waves with various wavelengths (and polarizations). It is a good approximation to take dispersion relation

Assume isotropic)  $\omega^2 = v^2 \vec{k}^2$  (no dispersion)

- where  $v$  is speed of sound

(Not exactly right at very short wavelength, but good enough)

- this is just like light (in radiation,



If the crystal is in a finite box (size  $L$ ), we'll have discrete values for  $\vec{k}$

$$\vec{k} = \frac{\pi}{L} \vec{n}$$

- Unlike light - there are 3 polarizations for each value of  $\vec{k}$  (longitudinal, too) (assume isotropy, and same  $v$  for all 3 pol.)

- Also unlike light - wavelength of light can be arbitrarily short. Not so for vibrations of crystal -- wavelengths less than interparticle spacing have no meaning.

Another way to say this - a system of  $N$  masses (in 3 dimensions) has  $3N$  degrees of freedom, and so there are  $3N$  normal modes (a finite number not scales with the crystal volume). To estimate maximum value of  $|k|$ :

4.15

$$\sum_n \rightarrow 3 \left(\frac{L}{\pi}\right)^3 \frac{1}{8} \int_{0}^{K_{\max}} 4\pi K^2 dK$$

$$(3 \text{ polarizations}) = \frac{3}{2} \frac{L^3}{\pi^2} \frac{L}{3} K_{\max}^3 = 3N$$

$$\boxed{\text{or } K_{\max}^3 = 6\pi^2 \frac{N}{V}}$$

called  $K_D$   
(per degree)

This makes sense

$$d_{\min}^3 = \left(\frac{2a}{K_{\max}}\right)^3 = \frac{8\pi^3}{6\pi^2} \frac{V}{N} = \frac{4}{3}\pi \frac{V}{N} = \frac{4}{3}\pi a^3$$

( $d_{\min}$  is actually  
volume of a sphere  
of radius  $a$ )  
But why?

Now — in quantum theory, energies of all  
the oscillators are quantized

$$\epsilon_\omega = n \hbar \omega$$

Sound is carried by quanta (analogous to photons)  
called phonons

A crystal at temperature  $\tau$  will  
vibrate. Let us calculate the energy stored  
in the vibrations. This is almost  
exactly the same as calculating the  
energy in a cavity. But with one two  
important differences — the wavelength  
cutoff (and the longitudinal polarization)

$$\text{So } U_V = \frac{3}{2\pi^2} \int_0^{K_D} dK K^2 \frac{\hbar \omega}{e^{\hbar \omega/\tau} - 1} \quad (3/2 \times \text{Planck})$$

And dispersion relation  $\omega = \nu k \Rightarrow$

$$U/V = \frac{3}{2\pi^2} \text{tr} \nu \left( \frac{\tau}{\text{tr} \nu} \right)^4 \int_0^{x_D} dx \frac{x^3}{e^x - 1}$$

$$X = \frac{\tau \nu k}{\tau} \text{ and } x_D = \frac{k \nu}{\tau} K_D = \Theta/\tau$$

(Define  $\Theta = \tau \nu k_D$  = Debye Temperature)

$$\Theta^3 = (\tau \nu)^3 \frac{6\pi^2 N}{V} \Rightarrow V = 6\pi^2 (\tau \nu)^3 \frac{N}{\Theta^3}$$

$$\Rightarrow U = 9 N \tau^4 \frac{1}{\Theta^3} \int_0^{x_D} dx \frac{x^3}{e^x - 1} \quad (\rightarrow f(\Theta/\tau))$$

Let's consider the limiting cases of large and small  $\tau$ :

(i) Low Temperature  $x_D = \Theta/\tau \gg 1$

Make an exponentially small error by extending integral to  $x = \infty$

$$I = \pi^4 / 15$$

$$\Rightarrow U = \frac{3\pi^4}{5} N \tau^4 \frac{1}{\Theta^3}$$

same  $\tau^4$  behavior  
as for B.B. radiation  
at low  $\tau$  - the  
wavelength cutoff

is unimportant because high  $\omega$  modes are unoccupied anyway

This formula works very well for solids at low temperature - the Debye  $T^3$  law

Typical values are  $\Theta \sim 200^\circ - 600^\circ K$

Note: called  $T^5$  law because

$$C_V = \left(\frac{\partial U}{\partial T}\right)_V = \frac{12\pi^4}{5} N \left(\frac{\tau}{\theta_B}\right)^3$$

- what one actually measures.

(4.17)

So  $\tau \lesssim 100^\circ K$  (liquid  $N_2$ ) is typically low enough.

(ii) High Temperature  $x_D = \theta/\tau \ll 1$

Then  $\int_0^{x_D} dx \frac{x^3}{e^x - 1} \sim \int_0^{x_D} dx x^2 = \frac{1}{3} x_D^3$

$$\boxed{U = 3N\tau} \quad \text{- this is just like Rayleigh-Jeans Law}$$

for phonons, integrated up to cutoff. It is a classical result (no  $\hbar$ )  $\Rightarrow$  i.e. equipartition: energy  $\propto$  in each oscillation

the high- $T$  limit  $\boxed{C_V = 3N}$

is known as Law of Du Mong and Pet.T., discovered (experimentally) in 1819.

And derived by Boltzmann, in 1876, by using classical statistical mechanics (equipartition)

It was Einstein (1906) who first understood that this result is  $T \rightarrow \infty$  limit, and that in quantum theory one should have

$C_V \rightarrow 0$  as  $T \rightarrow 0$  (first hint of 3rd law of thermodynamics)

Debye derived  $T^3$  law in 1912.

one of the great early experimental triumphs of quantum theory. (Einstein had assumed "monochromatic phonons" - all with the same  $\omega$  - instead of a frequency spectrum.)